

Janossy densities for chiral random matrices + application to 2c QCD

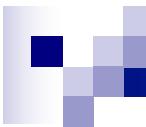
Shinsuke M. Nishigaki

Shimane Univ.

H. Fuji, I. Kanamori, SMN arXiv:1903.07176 (45 pp) = JHEP 09 (2019) ***

SMN arXiv:1606.00376 (LATTICE proceedings)

RMT in Sub-Atomic Physics and Beyond (in Honor of Jac Verbaarschot's 65th Birthday)
2019.8.5-9 ECT*, Trento, Italy



Janossy densities ^{§2}

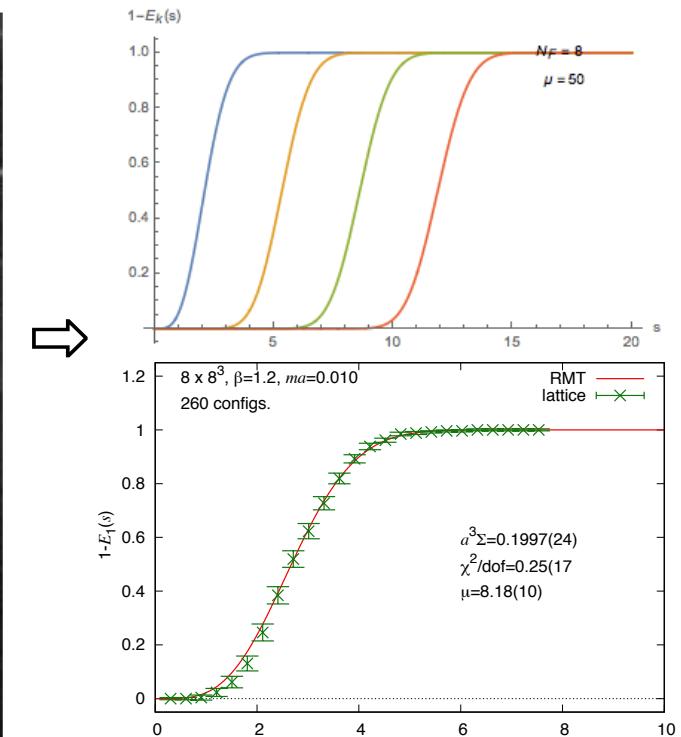
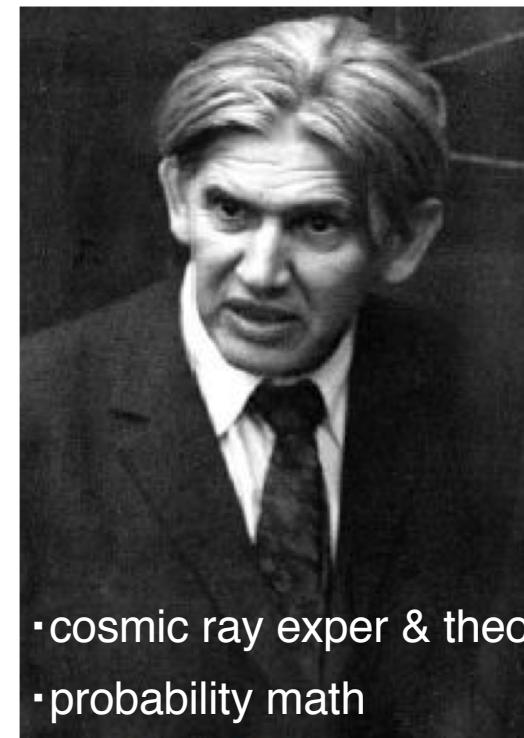
for chiral random matrices

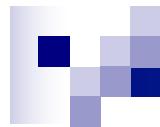
§4 Dirac spectrum of
+ application to 2c QCD

§1 Introduction

§3 Ordered EV
statistics

§5 Summary





personal recollection on Jac :

invited me & my wife to Thanksgiving dinner
at his house, for the 1st time - Nov. 2000



he cut & served a big turkey for us,
but didn't start eating it himself



S : why don't you eat it?

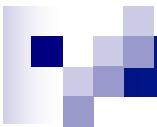
(Oh, he must have prepared it solely
to treat us, despite he's vegan…)

J : because I'm a little

worried about **Salmonella**



Happy Birthday, Jac



1. Introduction

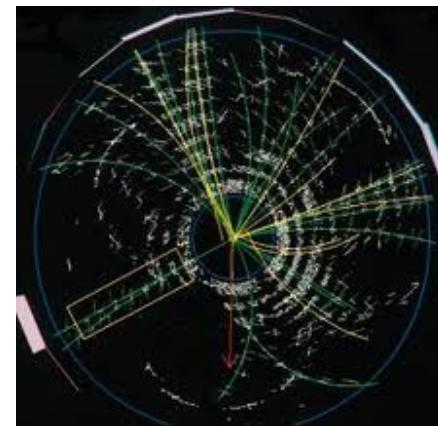
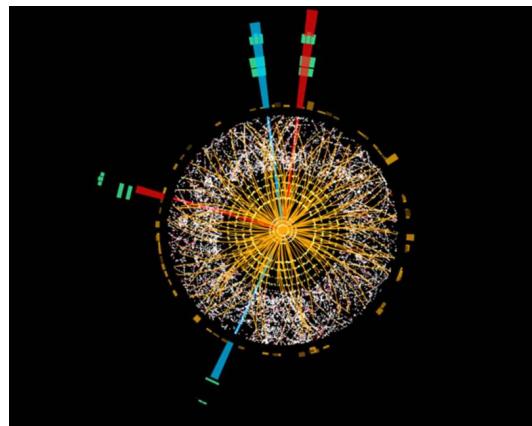
Hierarchy problem of SM

Higgs boson $M=125\text{GeV}$

LHC@CERN 2012;  2013

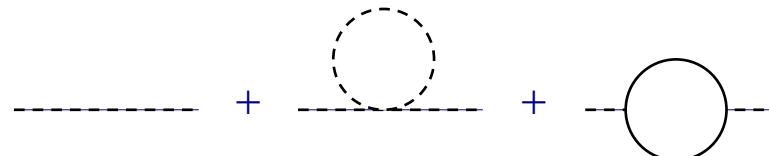
\approx Top quark $m=173\text{GeV}$

Tevatron@FNAL 1995;  2008



radiative correction of masses

$$M^2 = M_0^2 + \# \Lambda^2 + \# m_0^2 \log \frac{\Lambda}{m_0} + \dots$$



$$m = m_0 + \# m_0 \log \frac{\Lambda}{m_0} + \dots \quad \text{Weisskopf '34}$$



extremely-fine tuning needed to account for $M_H=125\text{GeV} \ll \Lambda \sim 10^{17}\text{GeV}$

Hierarchy problem of SM

Solution 1: SUSY

(H, ψ_H) degenerate $m_0 \Rightarrow$ degenerate m

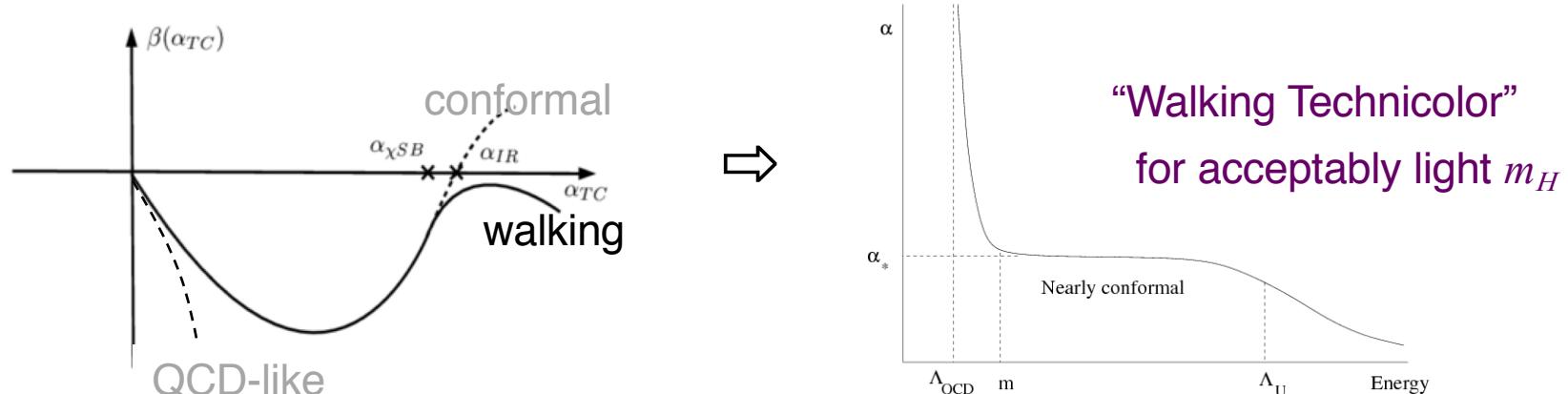
$$m^2 = m_0^2 + 0 + \# m_0^2 \log \frac{\Lambda}{m_0} + \dots$$

\uparrow
 $O(\Lambda^2)$ cancelled

$$m = m_0 + \# m_0 \log \frac{\Lambda}{m_0} + \dots$$

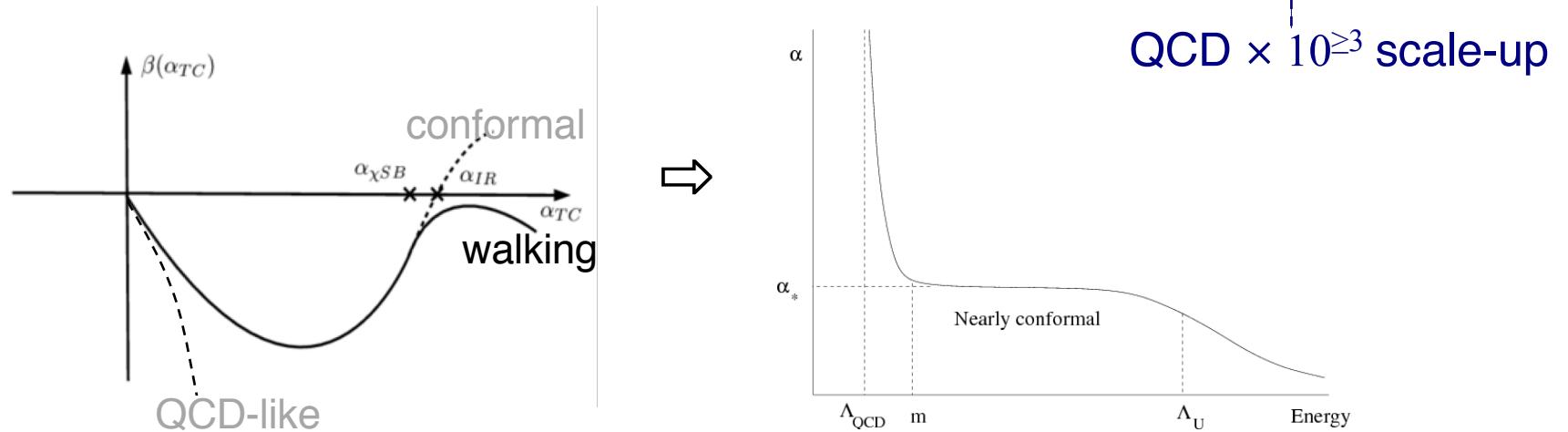
Solution 2: Higgs is fundamental \Rightarrow composite of 2 fermions in (new) gauge theory

$H \sim \bar{t} t$ (top condensation) or $\bar{Q} Q$ (technicolor)



Technicolor proposal

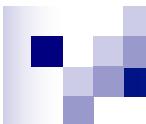
Solution 2: Higgs is fundamental \Rightarrow composite of 2 fermions in new gauge theory



$$\beta(\alpha) = - \underbrace{\left(\frac{11}{3}N - \frac{2}{3}n_F - \frac{4}{3}N n_{Ad} \right)}_{\text{small but } >0} \alpha^2 - \underbrace{\left(\frac{34}{3}N^2 - \left(\frac{13}{3}N - \frac{1}{N} \right)n_F - \frac{32}{3}N^2 n_{Ad} \right)}_{<0} \alpha^3 - \dots$$

- $N = 3, n_F = 12$ Nagoya KMI group; Kuti; Hasenfratz
- $N = 2, n_{Ad} = 2$ “Minimal WTC” Catterall-Sannino
- $N = 2, n_F = 8$ NCTU group, Helsinki group...

Chiral symm. broken like QCD, or
Conformal (unsuited for WTC) ?
...tested for SU(2) GTs on lattice



Chiral Random Matrices

[Shuryak-Verbaarschot '93]

$$\langle \dots \rangle_{\text{chRM}} = \frac{1}{Z(m)} \int dH e^{-\text{tr} H^2} \prod_{f=1}^{n_f} \det(H + i m_f) \dots , \quad H =$$

$$\left[\begin{array}{c|c} 0 & M \\ \hline M^+ & 0 \end{array} \right]^N$$

HS transf.

 $N \rightarrow \infty$, $m_f \rightarrow 0$ $\mu_f = N m_f$ fixed

captures all
global & discrete symmetries
of a gauge theory

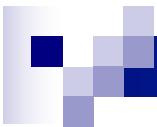
$$\int_{U(N_F)} dU (\det U)^v \exp\{\text{Re} \operatorname{tr} \operatorname{diag}(\mu_f) U\} = \text{cst.} \frac{\det [\mu_i^{j-1} I_{v+j-i}(\mu_i)]_{i,j=1}^{N_F}}{\Delta(\mu_f^2)}$$

↑
effective chL: $V \Sigma \operatorname{Re} \operatorname{tr} \operatorname{diag}(M_f) U$

$$M_{ab} \in \begin{cases} \mathbf{R} : \beta = 1 \\ \mathbf{C} : \beta = 2 \\ \mathbf{H} : \beta = 4 \end{cases}$$

if a gauge theory is in **chSB phase**,

- Dirac EVs $\{V\lambda_k\}$ measured at the same scaled quark masses $\{VM_f\}$ on different V collapse onto chRM results
- extrapolation to $M_f \rightarrow 0$ retains the best-fit $\Sigma \neq 0$

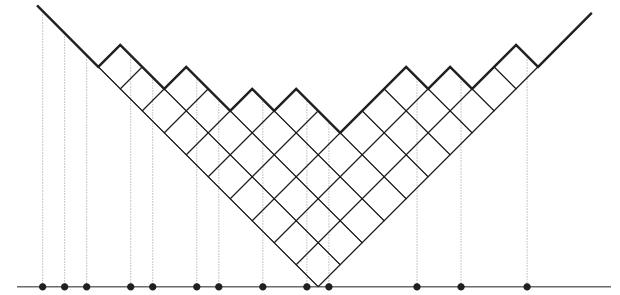


2. Janossy Density in DPP

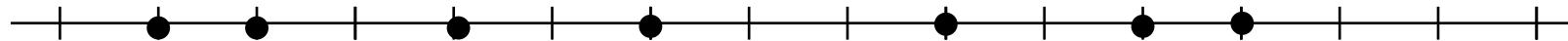
Determinantal point process

$$\text{Prob}_N(n_1, \dots, n_N) = \frac{1}{N!} \det \left[K(n_i, n_j) \right]_{i,j=1}^N \quad n_i \in \mathbf{Z}$$

$\mathbf{K} = [K(n,m)]_{n,m \in \mathbf{Z}}$: projective $\mathbf{K} \cdot \mathbf{K} = \mathbf{K}$, $\text{tr } \mathbf{K} = N$, $\mathbf{K}^\dagger = \mathbf{K}$



- Plancherel measure {YT}
- directed percolation
- continuous $\Rightarrow \beta=2$ RM



Janossy density

for Det point process $R_k(n_1, \dots, n_k) = \det[K(n_i, n_j)]_{i,j=1}^k$

$J_{k,I}(n_1, \dots, n_k)$ = Prob of



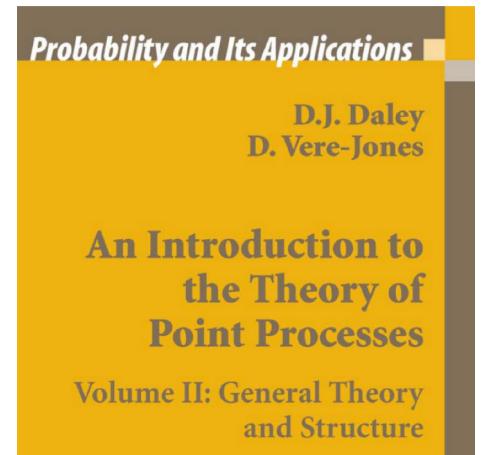
no pts in I except for k designated pts

[Daley-Vere-Jones '88, p.140]

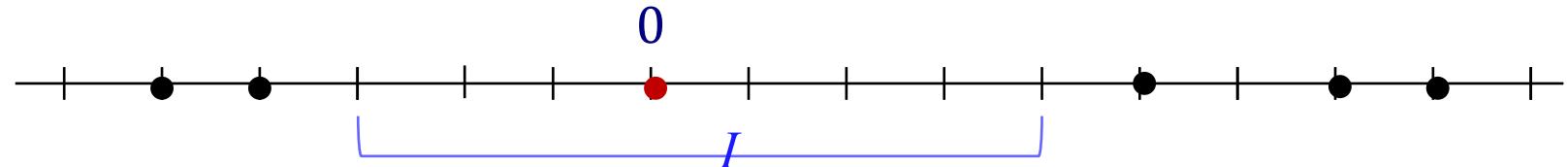
$$\mathbf{K}_I := [K(n,m)]_{n,m \in I}$$

$$J_{k,I}(n_1, \dots, n_k) = \det(1 - \mathbf{K}_I) \cdot \det \left[\left\langle n_i \middle| \mathbf{K}_I (1 - \mathbf{K}_I)^{-1} \middle| n_j \right\rangle \right]_{i,j=1}^k$$

all pts in I
designated pts ●



Sketch of Proof



$$J_{1,I}(0)$$

$$= R_1(0) - \sum_{n \in I} R_2(0, n) + \frac{1}{2!} \sum_{n, n' \in I} R_3(0, n, n') - \dots$$

$$= K(0, 0) - \sum_{n \in I} \begin{vmatrix} K(0, 0) & K(0, n) \\ K(n, 0) & K(n, n) \end{vmatrix} + \frac{1}{2!} \sum_{n, n' \in I} \begin{vmatrix} K(0, 0) & K(0, n) & K(0, n') \\ K(n, 0) & K(n, n) & K(n, n') \\ K(n', 0) & K(n', n) & K(n', n') \end{vmatrix} \dots$$

$$= \langle 0 | \mathbf{K}_I | 0 \rangle - \{ \langle 0 | \mathbf{K}_I | 0 \rangle \operatorname{tr} \mathbf{K}_I - \langle 0 | \mathbf{K}_I^2 | 0 \rangle \} \\ + \frac{1}{2!} \{ \langle 0 | \mathbf{K}_I | 0 \rangle (\operatorname{tr} \mathbf{K}_I)^2 - \langle 0 | \mathbf{K}_I | 0 \rangle \operatorname{tr} \mathbf{K}_I^2 - 2 \langle 0 | \mathbf{K}_I^2 | 0 \rangle \operatorname{tr} \mathbf{K}_I + 2 \langle 0 | \mathbf{K}_I^3 | 0 \rangle \} - \dots$$

$$= \langle 0 | \mathbf{K}_I | 0 \rangle \left\{ 1 - \operatorname{tr} \mathbf{K}_I + \frac{1}{2!} (\operatorname{tr} \mathbf{K}_I)^2 - \dots \right\} \left\{ 1 - \frac{1}{2} \operatorname{tr} \mathbf{K}_I^2 + \dots \right\} \dots$$

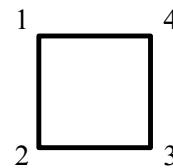
$$+ \langle 0 | \mathbf{K}_I^2 | 0 \rangle \{ 1 - \operatorname{tr} \mathbf{K}_I + \dots \} \dots$$

$$+ \langle 0 | \mathbf{K}_I^3 | 0 \rangle \{ 1 - \dots \} \dots + \dots$$

$$= \langle 0 | \mathbf{K}_I + \mathbf{K}_I^2 + \mathbf{K}_I^3 + \dots | 0 \rangle \det(\mathbf{1} - \mathbf{K}_I) = \langle 0 | \mathbf{K}_I (\mathbf{1} - \mathbf{K}_I)^{-1} | 0 \rangle \det(\mathbf{1} - \mathbf{K}_I)$$

2 Janossy density in DPP

Example: 2 Fermions on



$$\left. \begin{array}{l} \text{Prob}(\bullet \square) = 1/8 \\ \text{Prob}(\square \bullet) = 1/4 \end{array} \right\}$$

$$\Leftrightarrow \text{Prob}(n, m) = \frac{1}{2!} \begin{vmatrix} K(n, n) & K(n, m) \\ K(m, n) & K(m, m) \end{vmatrix}$$

$$\mathbf{K} = \frac{1}{4} \begin{pmatrix} 2 & 1-i & 0 & 1+i \\ 1+i & 2 & 1-i & 0 \\ 0 & 1+i & 2 & 1-i \\ 1-i & 0 & 1+i & 2 \end{pmatrix} = \mathbf{K} \cdot \mathbf{K}$$

$\text{tr } \mathbf{K} = 2$

$$J_{1,I}(1) = \text{Prob}(\boxed{I}) = \boxed{\bullet} + \boxed{\bullet} = \frac{1}{8} + \frac{1}{4}$$

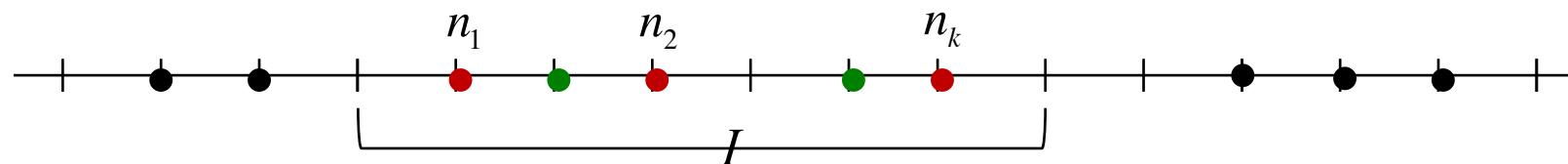
$$= \det(\mathbf{1} - \mathbf{K}_I) \cdot \langle 1 | \mathbf{K}_I (\mathbf{1} - \mathbf{K}_I)^{-1} | 1 \rangle$$

$$= \left| \mathbf{1} - \frac{1}{4} \begin{pmatrix} 2 & 1-i \\ 1+i & 2 \end{pmatrix} \right| \cdot \begin{pmatrix} 3 & 2-2i \\ 2+2i & 3 \end{pmatrix}_{1,1} = \frac{1}{8} \cdot 3$$

Janossy density

Det. point process: $R_k(n_1, \dots, n_k) = \det[K(n_i, n_j)]_{i,j=1}^k$

$J_{p,k,I}(n_1, \dots, n_k)$ = Prob of



exactly p pts \bullet in I except for k designated pts \bullet

$$\mathbf{K}_I := [K(n, m)]_{n, m \in I}$$

$$J_{p,k,I}(n_1, \dots, n_k) = \frac{1}{p!} (-\partial_\xi)^p \det(1 - \xi \mathbf{K}_I) \cdot \det \left[\langle n_i | \mathbf{K}_I (1 - \xi \mathbf{K}_I)^{-1} | n_j \rangle \right]_{i,j=1}^k \Big|_{\xi=1}$$

2 Janossy density in DPP

- Alternative expression by $\begin{vmatrix} A & B \\ C & D \end{vmatrix} = |D| |A - CD^{-1}B|$

$$J_{k,I}(n_1, \dots, n_k) = \det(\mathbf{1} - \mathbf{K}_I) \cdot \det[\langle n_i | \mathbf{K}_I (\mathbf{1} - \mathbf{K}_I)^{-1} | n_j \rangle]_{i,j=1}^k$$

$$= (-)^k \det \left| \begin{array}{c} \text{designated pts } \bullet \\ \hline -[\langle n_i | \mathbf{K}_I | n_j \rangle]_{i,j=1}^k \\ -[\langle n_i | \mathbf{K}_I | m \rangle]_{m \in I}^{j=1, \dots, k} \end{array} \right| \left| \begin{array}{c} \text{all pts in } I \\ \hline -[\langle m | \mathbf{K}_I | n_j \rangle]_{j=1, \dots, k}^{m \in I} \\ \mathbf{1} - \mathbf{K}_I \end{array} \right| \bullet_I$$

unifies 2 well-known formulas in DPP / RMT

$$R_k(n_1, \dots, n_k) = \det[K(n_i, n_j)]_{i,j=1}^k : k\text{-point Corr. Function}$$

$$E_{0,I} = \det(\mathbf{1} - \mathbf{K}_I) : \text{Gap Probability}$$

- Generalizable to qdet : Pfaffian PP = RMT at $\beta=1, 4$

2 Janossy density in DPP

- Alternative expression by $\begin{vmatrix} A & B \\ C & D \end{vmatrix} = |D| |A - CD^{-1}B|$

$$J_{k,I}(n_1, \dots, n_k) = \det(\mathbf{1} - \mathbf{K}_I) \cdot \det[\langle n_i | \mathbf{K}_I (\mathbf{1} - \mathbf{K}_I)^{-1} | n_j \rangle]_{i,j=1}^k$$

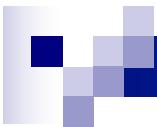
$$= (-)^k \det \begin{vmatrix} \underbrace{-[\langle n_i | \mathbf{K}_I | n_j \rangle]_{i,j=1}^k}_{\text{designated pts } \bullet} & \underbrace{-[\langle m | \mathbf{K}_I | n_j \rangle]_{j=1,\dots,k}^{m \in I}}_{\text{all pts in } I} \\ \underbrace{-[\langle n_i | \mathbf{K}_I | m \rangle]_{m \in I}^{j=1,\dots,k}}_{\mathbf{1} - \mathbf{K}_I} & \end{vmatrix} \bullet_I$$

- Continuous distribution : Fredholm Det, approx. by $I \sim \{y_1, \dots, y_m\}$

$$K(m, m') \rightarrow \sqrt{\Delta y_a} K(y_a, y_b) \sqrt{\Delta y_b}$$

$$J_{k,I}(x_1, \dots, x_k) = \lim_{\Delta y_a \rightarrow 0} (-)^k \det \begin{vmatrix} -[K(x_i, x_j)]_{i,j=1}^k & -[\sqrt{\Delta y_a} K(y_a, x_i)]_{y_a \in I}^{i=1,\dots,k} \\ -[K(x_i, y_b) \sqrt{\Delta y_b}]_{y_b \in I}^{j=1,\dots,k} & \mathbf{1} - [\sqrt{\Delta y_a} K(y_a, y_b) \sqrt{\Delta y_b}]_{y_a, y_b \in I} \end{vmatrix}$$

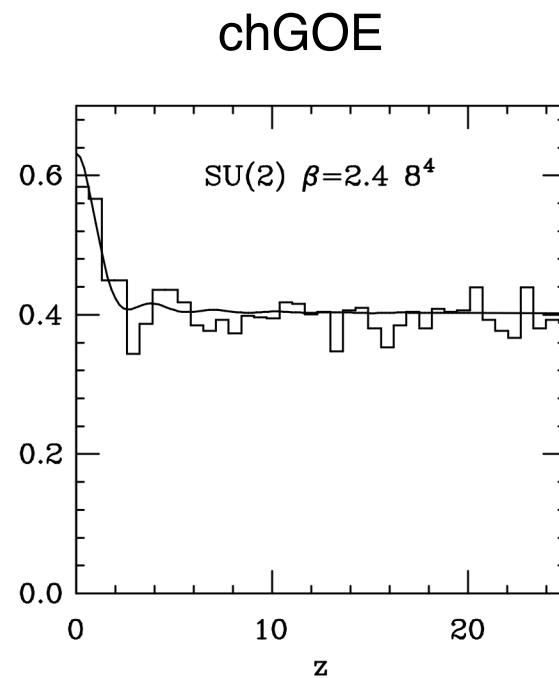
not explicit in [Borodin-Soshnikov '03][Forrester-Witte '07][Forrester-Witte-Bornemann '12]



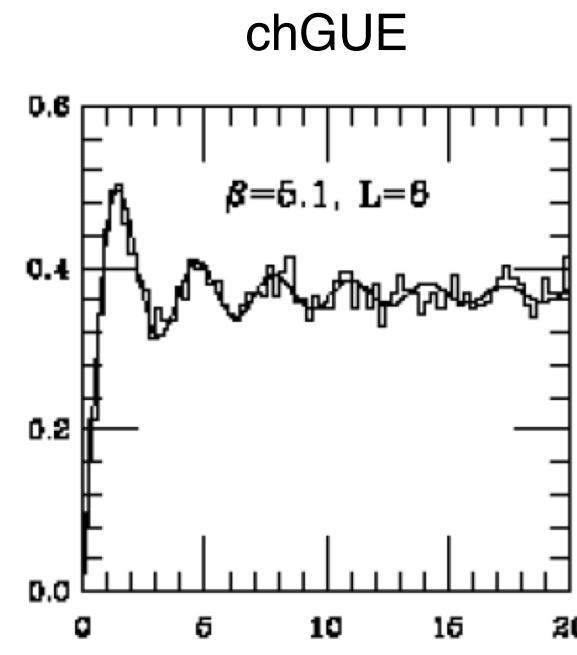
3. Ordered EV Statistics

Individual EV distributions

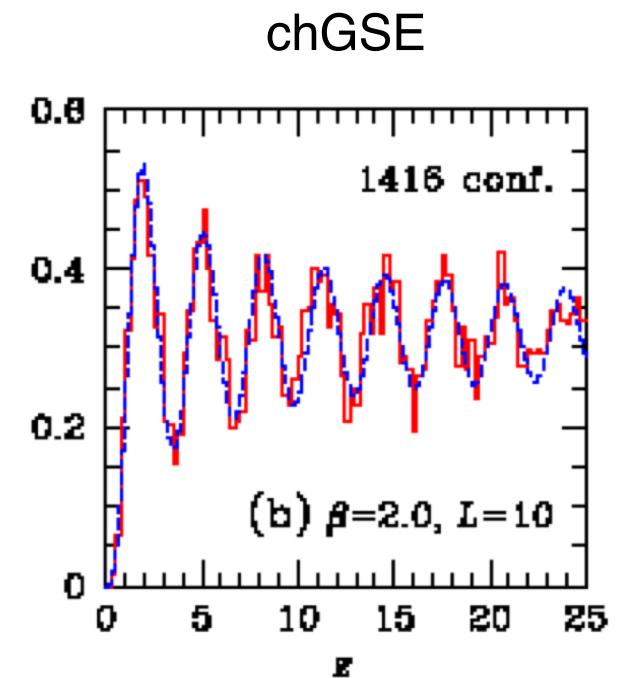
quotes from U. Heller's presentation: quenched chRMT vs LGT



[Edwards-Heller-Narayanan 99]



[Damgaard-Heller-Krasnitz 99]



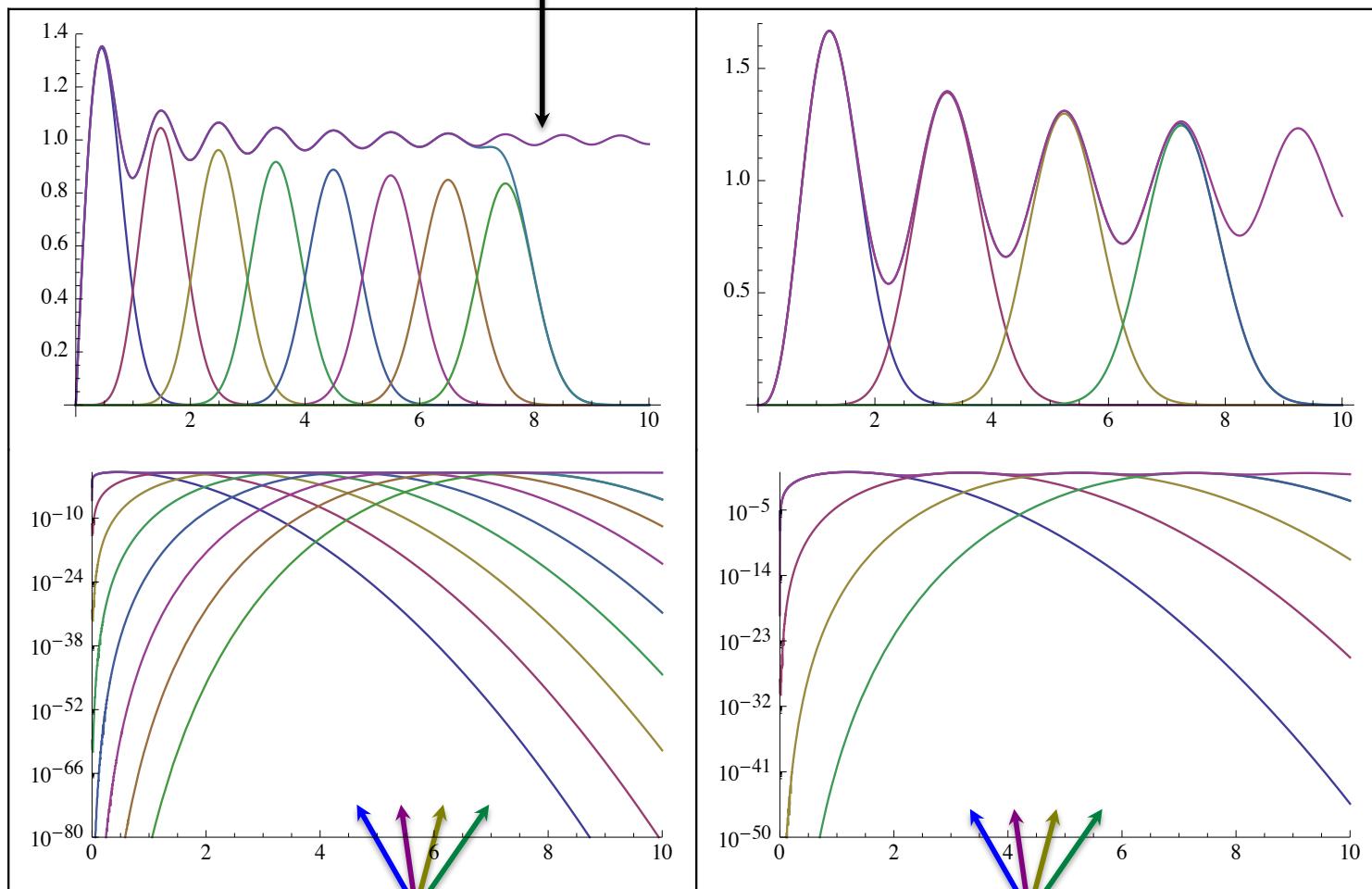
[Borbetti-Bitsch-Meyer-Schaefer-Verbaarschot-Wettig 98]

separation of peaks in $\text{Spec}(D)$ is highly desirable for precise fitting

Individual EV distributions

1st ~ 8th EV distributions of chGUE , chGSE

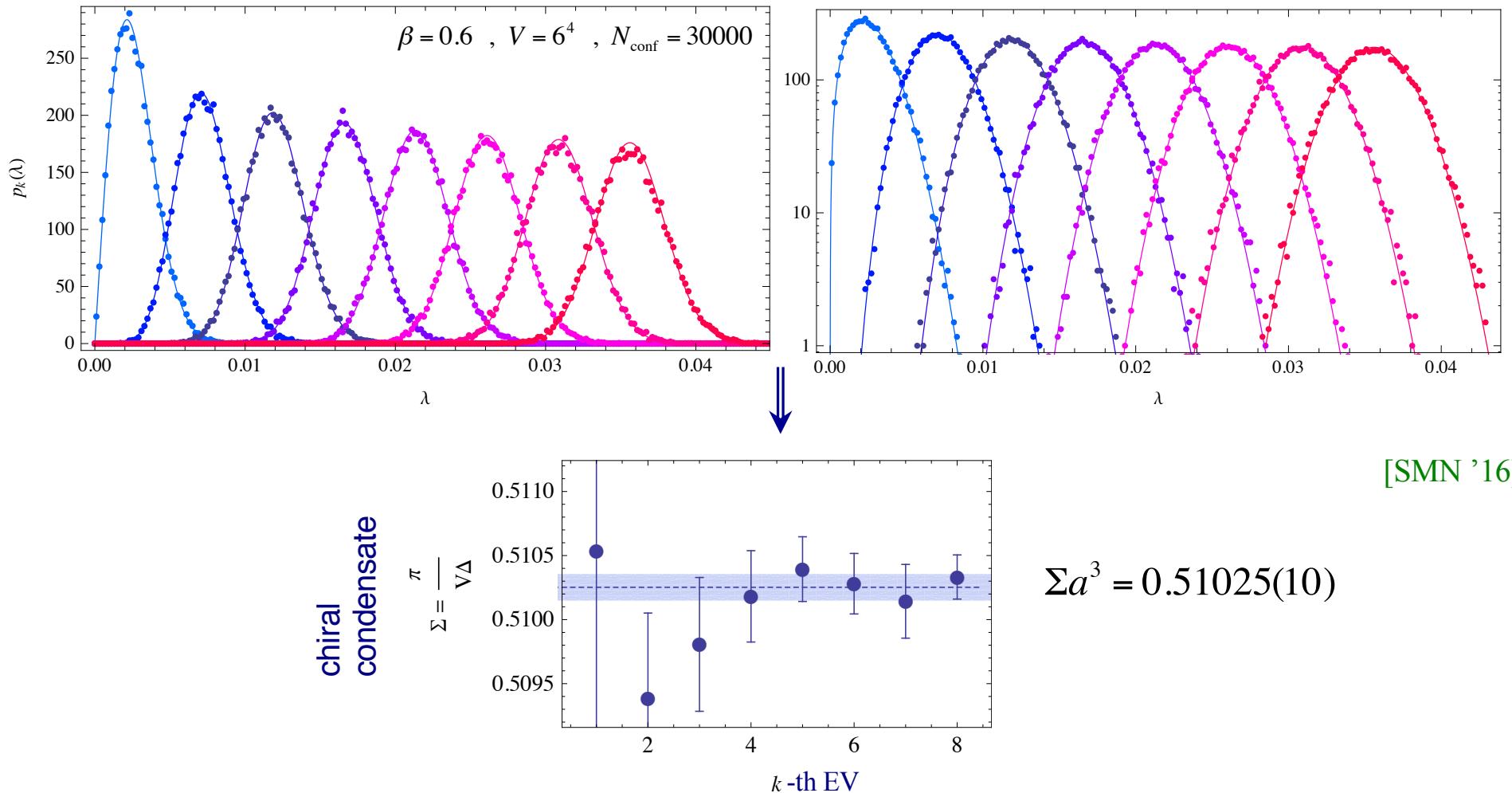
flattening oscillation \Rightarrow larger error for fitting



quasi-Gaussian \Rightarrow precise fitting possible

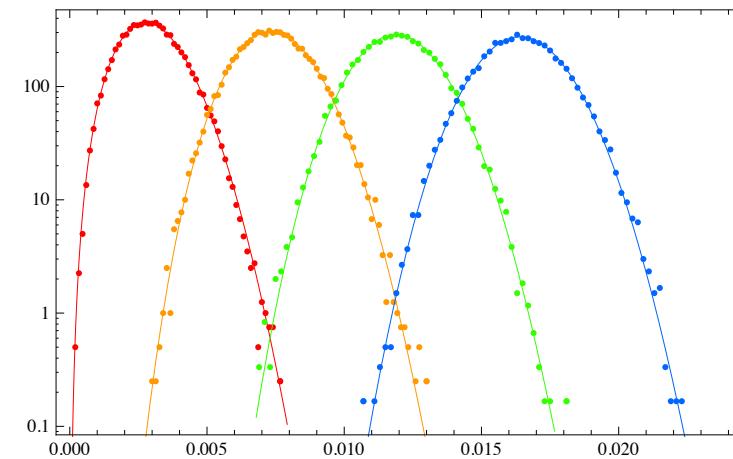
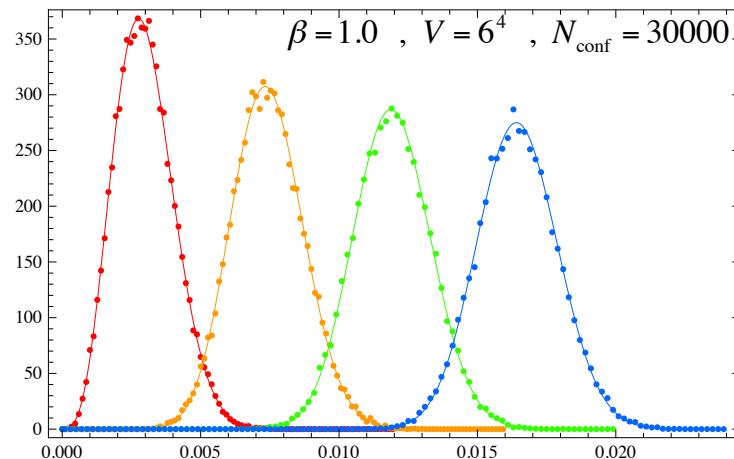
Chiral condensate from Individual EV distributions

exercise 1 : quenched U(1) Dirac spectrum vs chGUE

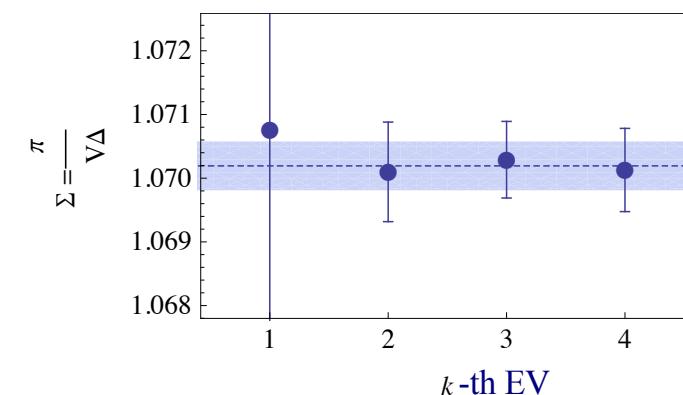


Chiral condensate from Individual EV distributions

exercise 2 : quenched SU(2) Dirac spectrum vs chGSE



chiral condensate



$$\Sigma a^3 = 1.07019(38)$$

[SMN '16]

Σ within $O(10^{-4})$ error!

Shifting method

[Damgaard-SMN '01]

$$\lambda_i \geq 0 \in \text{Spec}(H^2)$$

$$dH \ e^{-\text{tr} H^2} \prod_f \det(H + im_f) \propto \prod_{i=1}^N \left(d\lambda_i \ \lambda_i^{\beta(\nu+1)/2-1} \ e^{-\lambda_i} \prod_f (\lambda_i + m_f^2) \right) \prod_{i>j}^N |\lambda_i - \lambda_j|^\beta \quad \dots \text{ JPD of EVs}$$

$$P_k(\lambda_1, \dots, \lambda_k) = \boxed{\int_{\lambda_k}^{\infty} \cdots \int_{\lambda_k}^{\infty} d\lambda_{k+1} \cdots d\lambda_N \left(\text{JPD}_{N_F}(\lambda_1, \dots, \lambda_N; \{m\}; \nu) \right)} \quad \dots \text{ JPD of first } k \text{ EVs}$$

$\tilde{\lambda}_i \equiv \lambda_i - \lambda_k \quad \Downarrow$

$$= C(\{m\}) \int_0^{\infty} \cdots \int_0^{\infty} d\tilde{\lambda}_{k+1} \cdots d\tilde{\lambda}_N \left(\text{JPD}_{\tilde{N}_F}(\tilde{\lambda}_1, \dots, \tilde{\lambda}_N; \{\tilde{m}\}; \tilde{\nu}) \right) \xrightarrow{N \rightarrow \infty} \text{ratio of Bessel det's}$$

$$\Downarrow$$

$$p_k(\lambda_k) = \int_0^{\lambda_k} \cdots \int_0^{\lambda_k} d\lambda_1 \cdots d\lambda_{k-1} P_k(\lambda_1, \dots, \lambda_k) \xrightarrow{N \rightarrow \infty} \text{(k-1)-fold integral of ratio of Bessel det's}$$

for this trick to work, the exponent

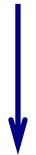
$$\beta \frac{\nu+1}{2} - 1 \in \begin{cases} \mathbf{N} & (\beta = 1, 2) \\ 2\mathbf{N} & (\beta = 4) \end{cases} \Rightarrow \begin{array}{l} \times \text{chGOE}, \ \nu = 0, 2, 4, \dots \\ \times \text{chGSE}, \ N_F = 0, 2, 4, \dots \end{array}$$

practically unfeasible
for $k \geq 5$

authors	year	β	k	ν	dynamical n_F	cons
Edelman	88, 91	1	1	0~3	0	finite N
Forrester	93	2	1	A	0	Fredholm det
		4	1	A	0	Jack poly.
Tracy-Widom	94	2	A	A	0	Painleve
Forrester-Hughes	94	2	1+2	A	0	shifting
Nagao-Forrester	98	4	1	A	0+odd massless	skew OP
Wilke-Guhr-Wettig	98	2	1	A	A	shifting
Berbenni-Bitsch-Meyer-Wettig	98	4	1	0	0 +evn massless	finite N
		4	1	0	1	finite N
Nagao-SMN, Damgaard-SMN	00, 01	2	A	A	A	shifting
		4	A	A	\forall pairs +odd	massless
		1	A	odd	A	
Forrester	00, 06	4, 1	A	A	0	Painleve
Wirtz-Akemann-Guhr-Kieburg-Wegner	14	1	1	evn	0	skew OP
SMN	16	2	A	A	A	Nystrom
Fuji-Kanamori-SMN	19	4	A	A	\forall pairs	Nystrom

Nystrom-type approx to Fredholm Det

Gauss-Legendre Quadrature : $\{x_1, \dots, x_M\} \in I, \quad \{\Delta x_1, \dots, \Delta x_M\} > 0$



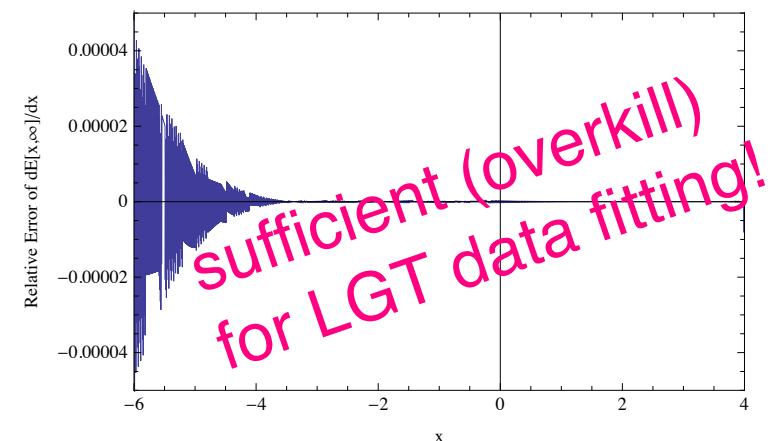
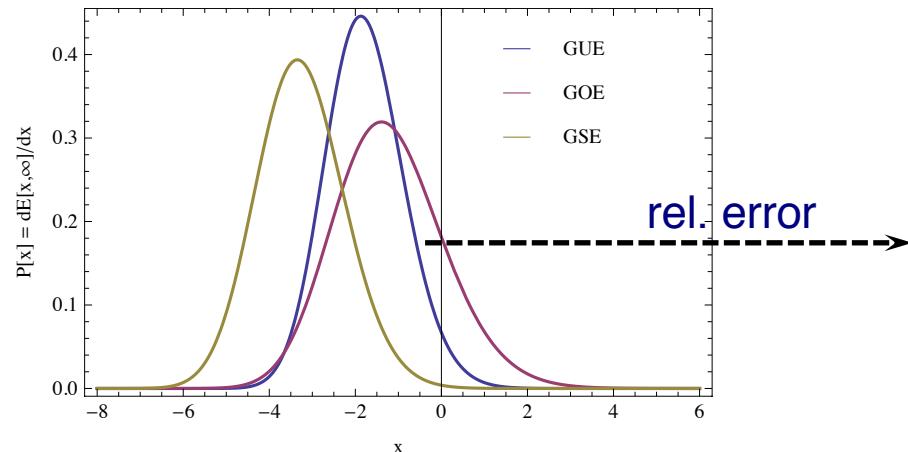
$$\int_I f(x) dx \cong \sum_{i=1}^M f(x_i) \Delta x_i, \text{ exact for } f(x) = x^M + \text{lower}$$

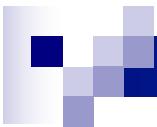
$$\text{Det}(1 - K_I) \cong \det \left[\delta_{ij} - K(x_i, x_j) \sqrt{\Delta x_i \Delta x_j} \right]_{i,j=1}^M + \text{relative error } O(e^{-\text{const.} M})$$

[Bornemann '10]

ex. Largest EV distribution

Nystrom approx ($M=30$) for K_{Airy} vs Tracy-Widom's analytic formula



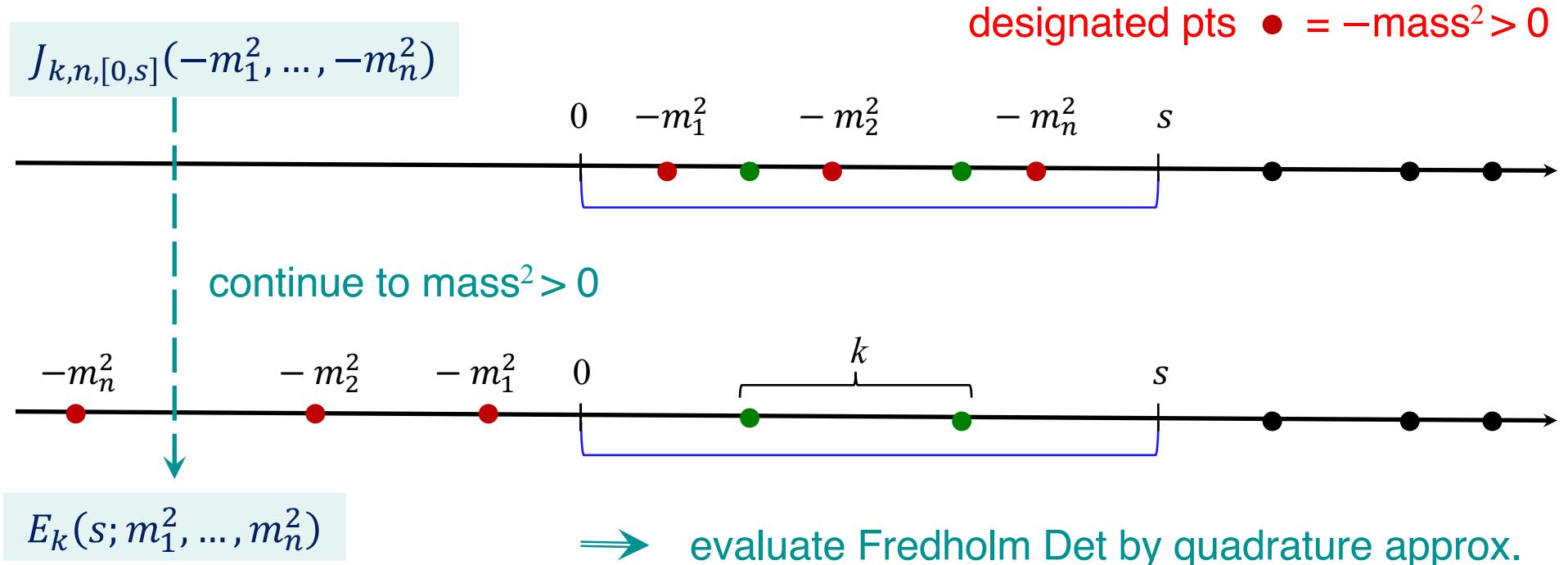


4. Dirac spectra of 2-color QCD

$n_F = \beta n$ fermions as Janossy density

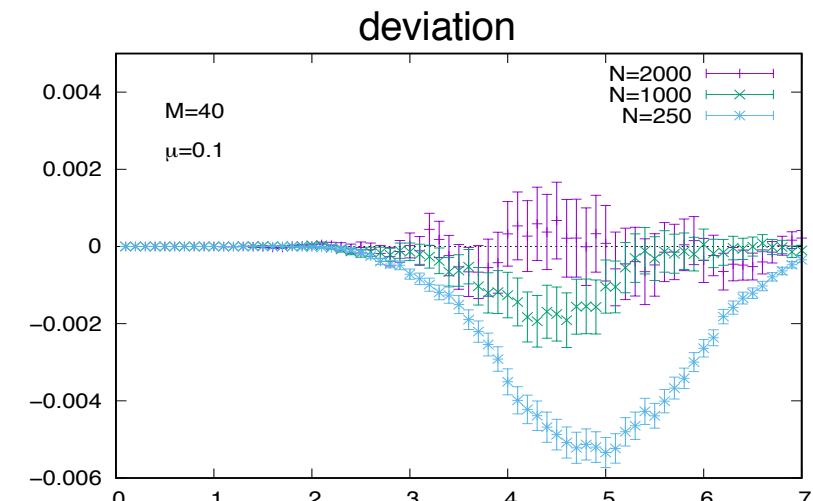
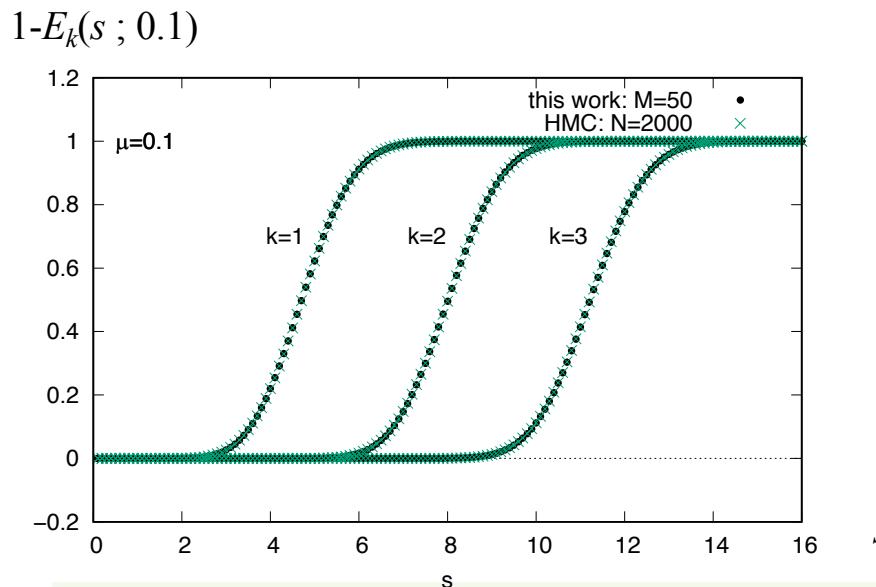
$$\int dH e^{-\text{tr} H^2} \prod_{f=1}^n (H + im_f)^\beta = \int_0^\infty \prod_{i=1}^N \left(d\lambda_i \lambda_i^{\frac{\beta(\nu+1)}{2}-1} e^{-\lambda_i} \prod_{f=1}^n |\lambda_i + m_f^2|^\beta \right) \prod_{i>j}^N |\lambda_i - \lambda_j|^\beta$$

$$\propto \int_0^\infty \prod_{i=1}^{N+n} \left(d\lambda_i \lambda_i^{\frac{\beta(\nu+1)}{2}-1} e^{-\lambda_i} \right) \prod_{i>j}^{N+n} |\lambda_i - \lambda_j|^\beta \cdot \prod_{k=N+1}^{N+n} \delta(\lambda_k - (-m_f^2))$$



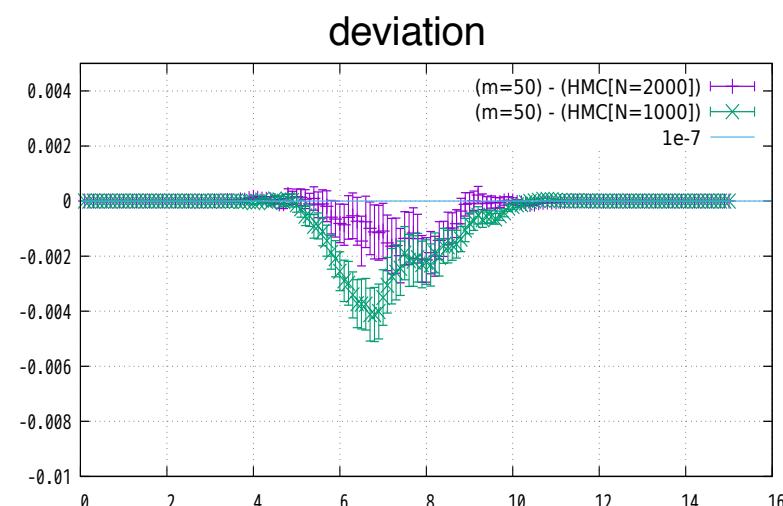
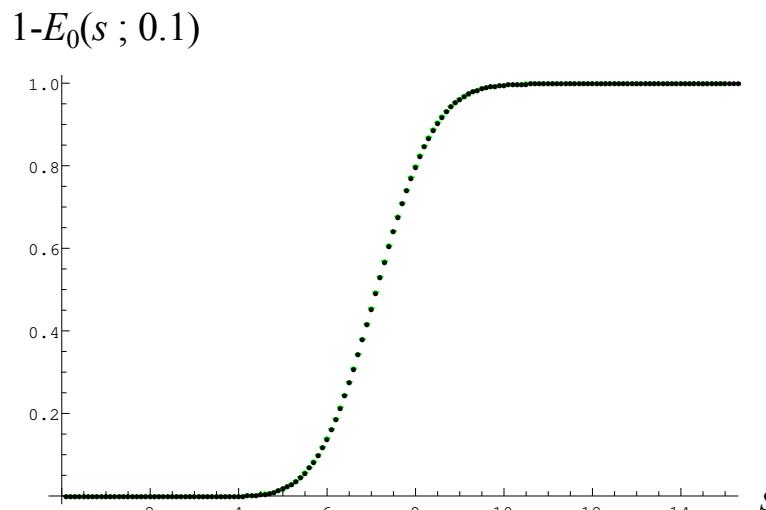
4 Dirac Spectra of 2C QCD

$\beta=4$ (chGSE), $n_F=4$, $k=0, 1, 2$, $\mu=0.1$ vs chGSE ($N=250\sim 2000$) by HMC



Quadrature Approx of Det : systematic error too small for bare eyes!

$\beta=4$ (chGSE), $n_F=8$, $k=0$, $\mu=0.1$ vs chGSE ($N=1000\sim 2000$) by HMC

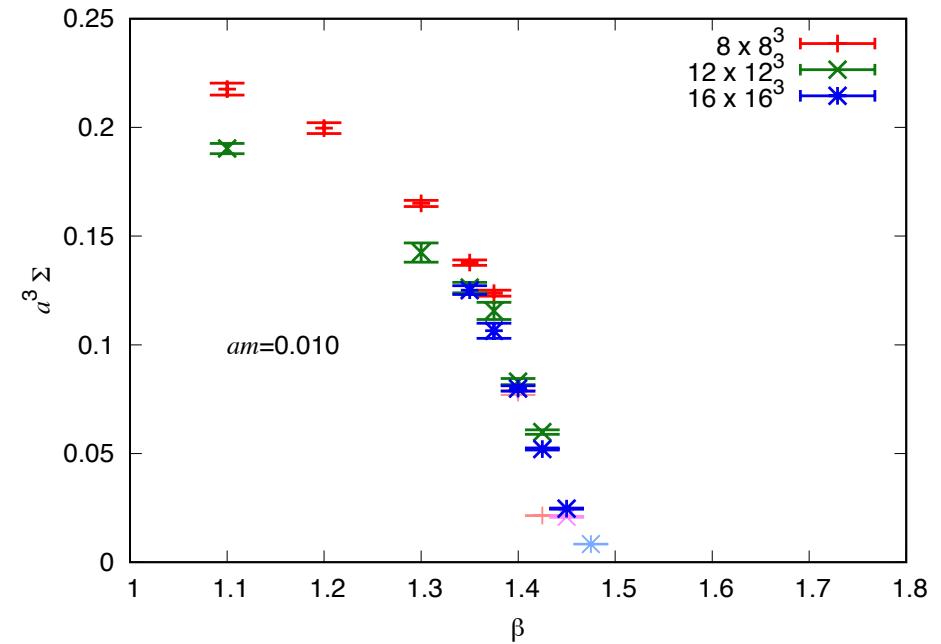
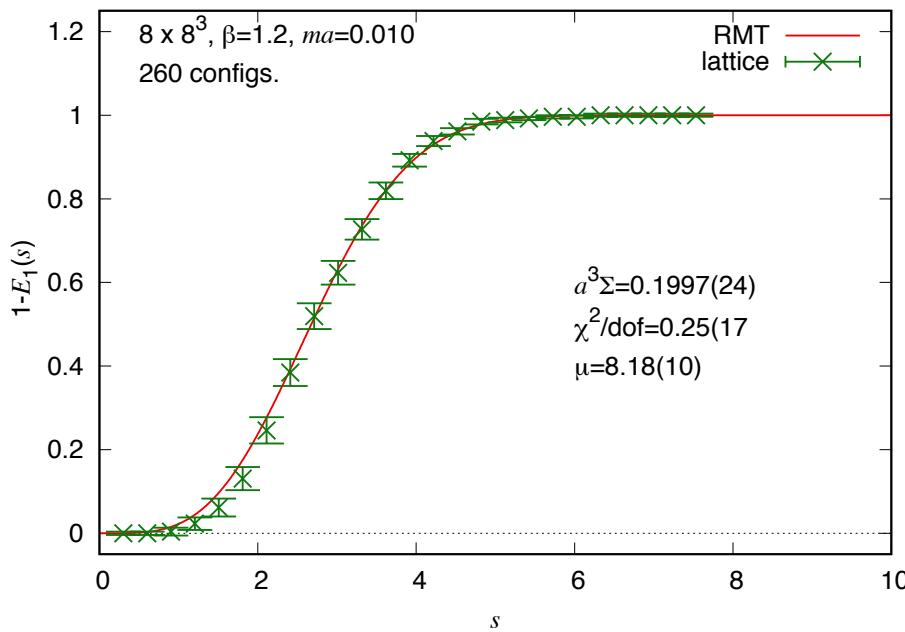


chSB check for WTC candidate

$\beta = 4$ (chGSE), $n_F = 4$, $k = 0$
 $\mu = 8.18$

vs
 2c QCD, $n_F = 2 \times 4 = 2 + (2+2+2)$
 $ma = 0.010$

flavor taste

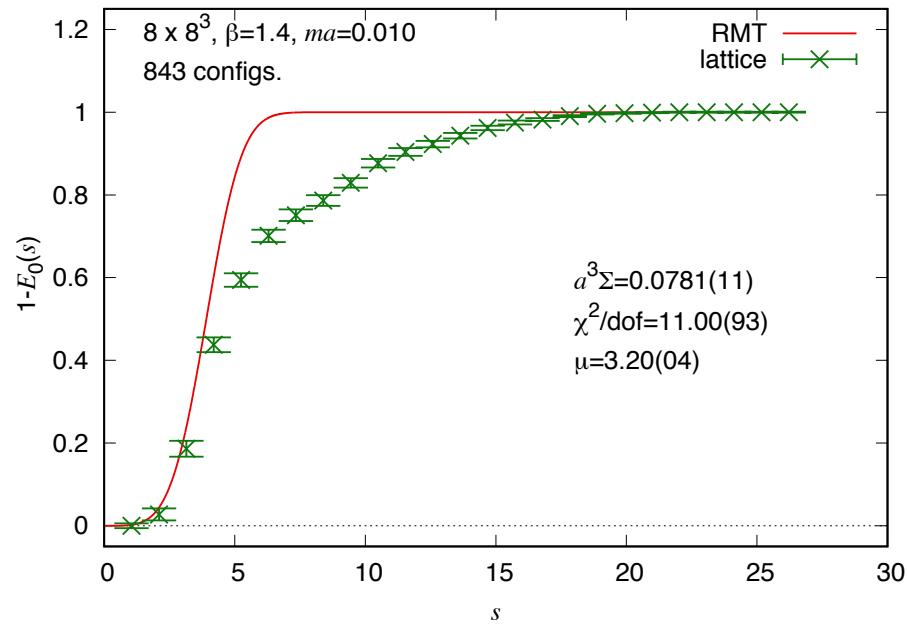


$N=2$, $n_F = 2 + (2+2+2)$, $\beta \leq 1.4\dots$: Chiral Symm Broken

chSB check for WTC candidate

$\beta = 4$ (chGSE), $n_F = 4$, $k = 0$

$\mu = 3.20$

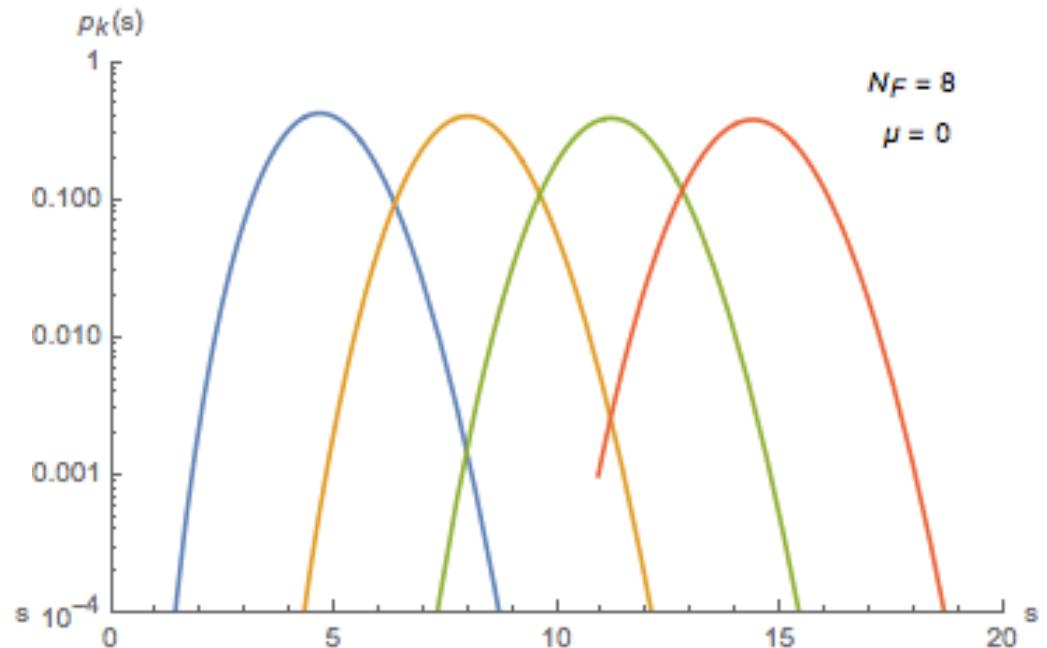
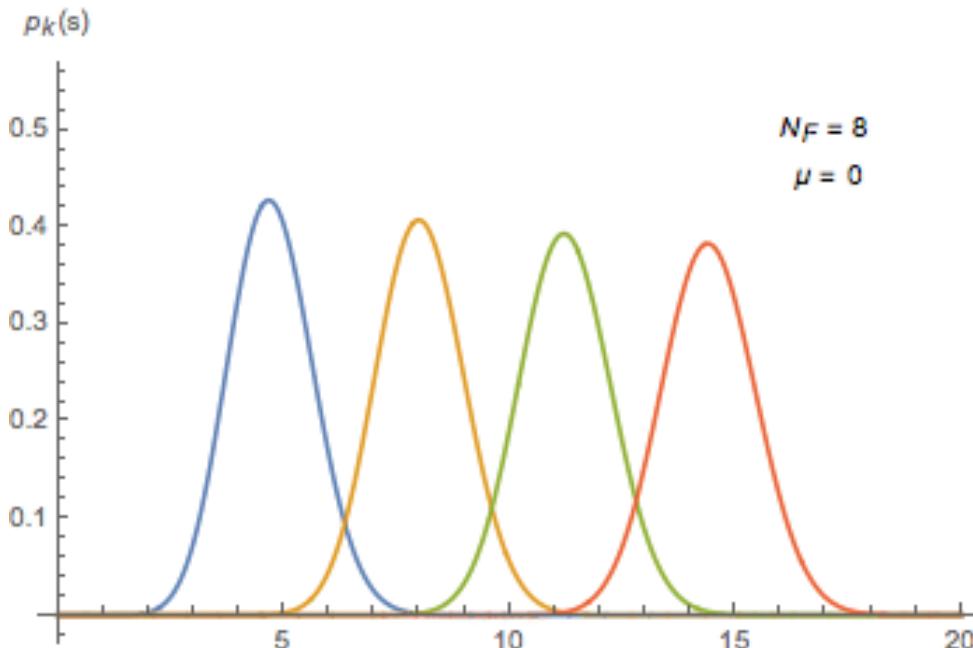


... bad χ^2 for chSym phase

$N=2$, $n_F = 2+(2+2+2)$, $\beta \leq 1.4\dots$: Chiral Symm Broken

chSB check for WTC candidate

$\beta = 4$ (chGSE), $n_F = 8$, $k = 0, 1, 2, 3$, $\mu = 0, 1, 2, \dots, \infty$



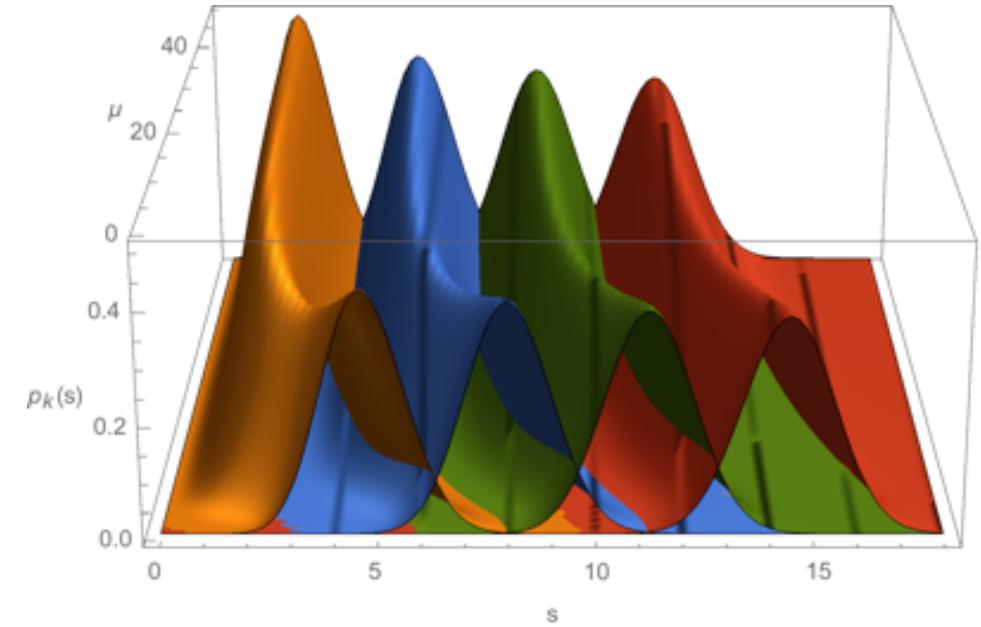
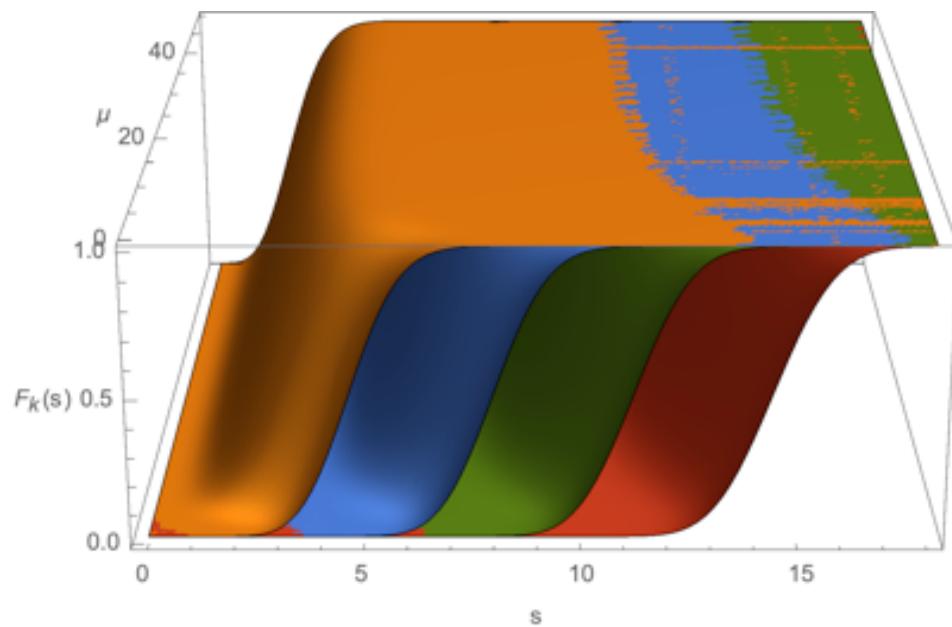
to be fitted with

flavor taste

$N = 2$, $n_F = 8 \times 1$ lattice QCD to judge whether chSB or conformal

chSB check for WTC candidate

$$\beta = 4 \text{ (chGSE)}, \ n_F = 8, \ k = 0, 1, 2, 3, \ \mu = 0, 1, 2, \dots, \infty$$



to be fitted with

flavor taste

$N = 2, \ n_F = 8 \times 1$ lattice QCD to judge whether chSB or conformal



§5 Summary

- 2 technical difficulties in evaluating Individual EVDs of massive chRMMs are overcome by Janossy Density formula + Quadrature method
- Individual Dirac EVDs of 2C QCD with $n_F = 4n$ staggered quarks, if the theory in χ SB phase, are predicted from massive chGSE
- Chiral cond Σ of 2C QCD with $n_F = 2 + (2+2+2)$ is determined by fitting $\text{Spec}(D)$ judged to be in the chSB phase (caution: $\beta = 4/g^2 \leq 1.3$)

- feasible plan : determine whether the WTC candidates in sympl. class 2C QCD with $n_F = 8$ (stag.) , $n_{Ad} = 2$ (overlap) is chSB, conformal, or else